

The Prospects for Quantum Information Processing Using Electron Spin in Semiconductors

A Semiconductor Based Teleportation Repeater For Long Distance Quantum Communications:

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1 qu-bit \equiv 1 spin

1 spin: $|\Psi\rangle = c_1|0\rangle + c_2|1\rangle$

2 spins: $|\Psi\rangle = c_1|00\rangle + c_2|01\rangle + c_3|10\rangle + c_4|11\rangle$

3 spins: $|\Psi\rangle = c_1|000\rangle + c_2|001\rangle + c_3|010\rangle + c_4|011\rangle + c_5|100\rangle + c_6|101\rangle + c_7|110\rangle + c_8|111\rangle$

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1 spin \rightarrow 2 complex numbers

2 spins \rightarrow 4 complex numbers

3 spins \rightarrow 8 complex numbers

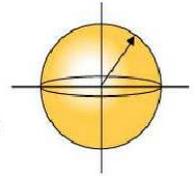
N spins $\rightarrow 2^N$ complex numbers

No. of digital classical configurations = 2^N

No. of unentangled quantum configurations = $(2^\alpha)^N$

($\alpha \equiv$ bits of analog precision)

No. of entangled quantum configurations = $(2^\alpha)^{2^N}$



2^{2^N} scaling:

How many classical bits do we need to specify a wave function with 2 bit accuracy?

Consider N=34 atoms

$$2^{34} \times 2^2 = 2^{36} = 80 \text{ Gbits} = 10 \text{ Gbytes}$$

“1 hard disk”

Consider N=61 atoms

$$2^{61} \times 2^2 = 2^{63} = 8 \times 10^{18} \text{ bits} = 10^8 \text{ hard disks}$$

“world-wide annual production”

Consider N=94 atoms

$$2^{94} \times 2^2 = 2^{96} = 8 \times 10^{28} \text{ bits} = 10^{18} \text{ hard disks}$$

“world-wide continuous production for the age of the universe”

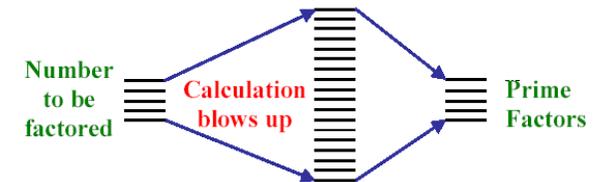
ordinary bit: $\alpha|0\rangle + \beta|1\rangle$

Now allow for conditional probability (also called entanglement)

$$\alpha|0\rangle + \beta|1\rangle$$

α, β depend on what all the other (N-1) spins are doing
 $\Rightarrow 2^{N-1}$ conditional probabilities and Huge Complexity

- You do not have internal access to quantum complexity.
- The entanglement collapses when you measure.
- There are N output bits and only 2^N possible answers.



You do not have access to the intermediate computations

Algorithms:

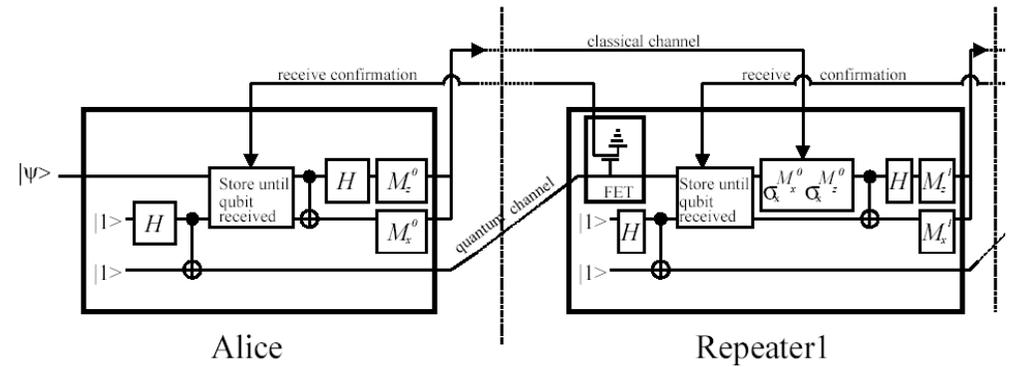
Shor's factorization algorithm (1994)

Kitaev's factorization algorithm (1996)

Near term:

Small computer for error correction of Quantum Cryptography-

"Quantum Repeater"



Error Correction:

Use redundancy, substitute 3 spins for one:
 $|1\rangle \Rightarrow |1\rangle|1\rangle|1\rangle$ Robust against spin-flips
 $|0\rangle \Rightarrow |0\rangle|0\rangle|0\rangle$

Phase Error Correcting Code:
 $|a\rangle \Rightarrow (|0\rangle + |1\rangle)(|0\rangle + |1\rangle)(|0\rangle + |1\rangle)$
 $|b\rangle \Rightarrow (|0\rangle - |1\rangle)(|0\rangle - |1\rangle)(|0\rangle - |1\rangle)$
 $|a\rangle \Rightarrow |000\rangle + |001\rangle + |010\rangle + |011\rangle + |100\rangle + |101\rangle + |110\rangle + |111\rangle$
 $|b\rangle \Rightarrow |000\rangle - |001\rangle - |010\rangle + |011\rangle - |100\rangle + |101\rangle + |110\rangle - |111\rangle$

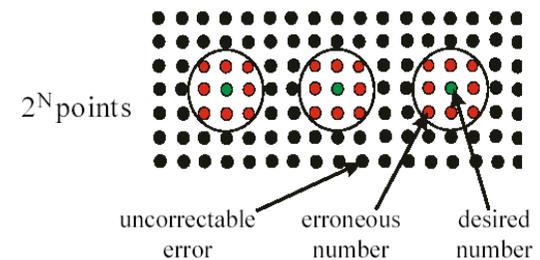
More compact phase error correcting code:
 $|a\rangle \Rightarrow |000\rangle + |011\rangle + |101\rangle + |110\rangle$
 $|b\rangle \Rightarrow |001\rangle + |010\rangle + |100\rangle + |111\rangle$

Shor's original phase and amplitude error correcting code:
 $|a\rangle \Rightarrow (|000\rangle + |111\rangle)(|000\rangle + |111\rangle)(|000\rangle + |111\rangle)$
 $|b\rangle \Rightarrow (|000\rangle - |111\rangle)(|000\rangle - |111\rangle)(|000\rangle - |111\rangle)$
 requires 9 qubits.

Threshold:
 Error Rate must be $< 10^{-3}$ or 10^{-4}

Classical Error Correction:

N bits plotted in N-dimensional space:



Because N-dimensional space is huge, the overhead on un-used points is relatively minor.

(N-k) dimensions of useful numbers.
 k dimensions of redundancy.

- Floppy disks
- CD-ROMs
- Digital Communication
- Voice Communication

Simplest case of 3-dim space,
 $2^3 = 8$ points

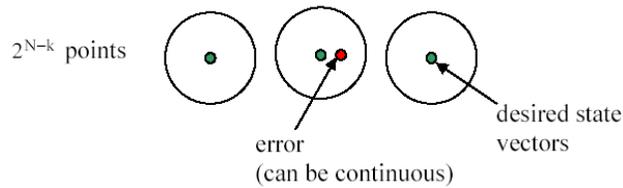
Code Errors

$|000\rangle$ $\left\{ \begin{array}{l} |100\rangle \\ |010\rangle \\ |001\rangle \end{array} \right.$

$|111\rangle$ $\left\{ \begin{array}{l} |011\rangle \\ |101\rangle \\ |110\rangle \end{array} \right.$

Quantum Error Correction:

(N-k) dimensional Hilbert space



Unitary rotation of error into the k redundant dimensions.

Errors build up in the redundant degrees of freedom, and must be purged from time to time.

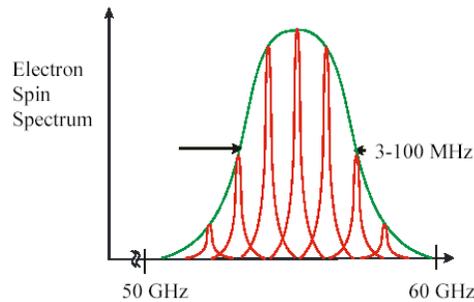
Quantum Error Correction works, only if errors are $<10^{-4}$

Physical realizations of qubits:

- photon polarization
- electron spin
- nuclear spin
- electron states in a trapped ion/atom
- Anharmonic LC oscillator states in a Josephson junction ring

Spins are favorable because of:

- Less environmental disturbance
- Less spontaneous emission and quantum noise at RF frequencies.
- electron spin transitions respond faster and easier to detect than nuclear spins
- But, superconducting rings are limited mainly by dissipation, that can be low



Electron spin in Si at $\sim 1^\circ\text{K}$ $\left\{ \begin{array}{l} T_1 \geq 1 \text{ hour} \\ 10 \text{ nsec} \leq T_2 \leq 500 \text{ nsec} \\ T_2 \geq 0.5 \text{ msec} \end{array} \right.$

Nuclear Spin $T_1 > 1 \text{ hour}$.
 Nuclear Spin T_2 is not measured.

Spin-Echo measurements:

J.P. Gordon et al, Physical Review Letters 1, 368 (1958)

M. Chiba & A. Hirai, J. Phy. Society of Japan 33, 730 (1972)

Electron Spin
 Line Broadening Mechanisms:

1. Spin-orbit coupling
 vibrations change the g-factor;
 proportional to T^4

2. Dyakonov $\vec{v} \times \vec{E}_{\text{eff}}$ mechanism (in III-V's)

3. Hyperfine coupling to nuclei:

$$\sum_n \vec{\mu}_e \cdot \vec{\mu}_n \frac{8\pi}{3} |\Psi(r_n)|^2 + \frac{3(\vec{\mu}_e \cdot \hat{r})(\vec{\mu}_n \cdot \hat{r}) - \vec{\mu}_e \cdot \vec{\mu}_n}{r^3}$$

isotopically pure ^{28}Si or Ge: 99.92% ^{28}Si is available.

4. Electrical Noise on the gates

5. Electron Spin-Electron Spin interaction;
 Figure-of-Merit = On/Off ratio of Exchange Interaction

In the chart of nuclides there are no spin zero nuclei in column III or column V of the periodic table!

III	IV	V
5 B ¹¹	6 C ¹² 99% spinless	7 N ¹⁴
13 Al ²⁷	14 Si ²⁸ 95% spinless	15 P ³¹
31 Ga ⁶⁹	32 Ge ⁷⁴ 92% spinless	33 As ⁷⁵
49 In ¹¹⁵	50 Sn ¹²⁰	51 Sb ¹²¹

What do we really know about T_2 & T_2^* ?

Electron Spins in Silicon at low temps			Electron Spins in III-V's at low temps		
	Homogeneous dephasing time T_2	Inhomogeneous dephasing time T_2^*		Homogeneous dephasing time T_2	Inhomogeneous dephasing time T_2^*
Trapped electrons	0.3msec-0.5msec <small>Spin Echo in n-Si</small>	100nsec-1 μ sec <small>ESR in n-Si</small>	Trapped electrons	Not Known! <small>Spin Echo in very lightly doped n-GaAs, to avoid exchange interactions</small>	≈ 3 nsec <small>ESR in n-GaAs</small>
Mobile electrons	100nsec-10 μ sec <small>magneto-resistive ESR in 2d Si-Ge electron-gas</small>	No Inhomogeneous Broadening	Mobile Electrons	0.1nsec-500nsec <small>(a) Awschalom et al, optical ESR (b) magneto-resistive ESR in 2-d e-gas</small>	No Inhomogeneous Broadening

Table 1. A short list of "long" T_2 values in comparison with the value found in n-type Si:

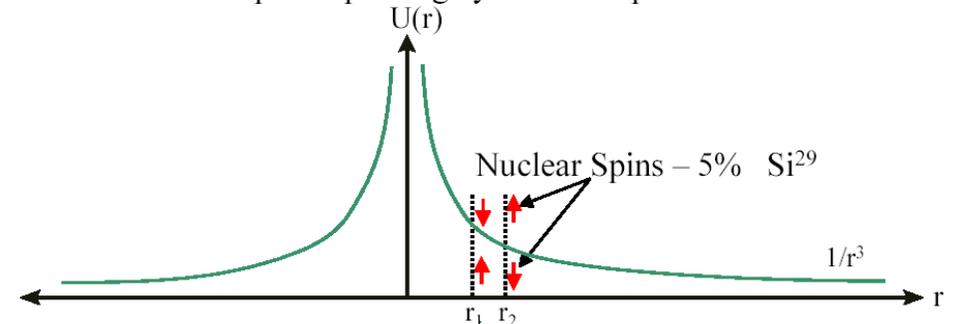
Material	temperature	T_2
γ -irradiated quartz	300K	2-5 μ sec
N-V centers in diamond, (price/purity dependent)	300K	2-100 μ sec
N@C60 (spin 3/2) ¹	300K solution	120 μ sec
Phosphorus in Si	<12°K	250 μ sec

Benefits of a Carbon Host for electron spin:

- Low Z \Rightarrow small spin-orbit dephasing.
 - High Debye temperature \Rightarrow less phonon dephasing.
 - 99% spin free C¹², only 1% C¹³, less hyperfine broadening.
 - Larger binding energies and energy spacings
- 300°K quantum information processor?

¹Knapp et al, Chem. Phys. Lett. 272, p.433 (1997)

Electron Spin Dephasing by Nuclear Spin Diffusion

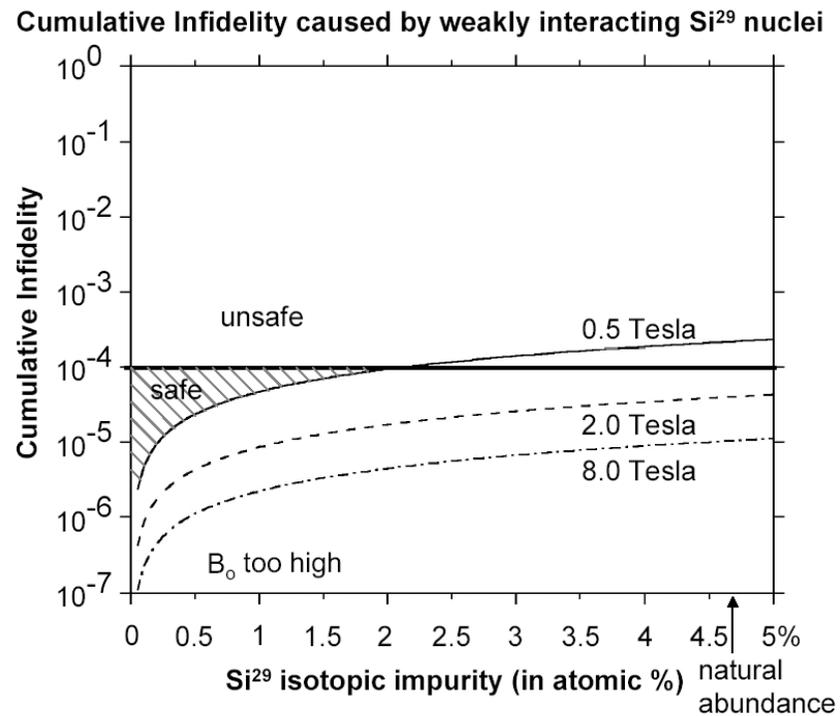
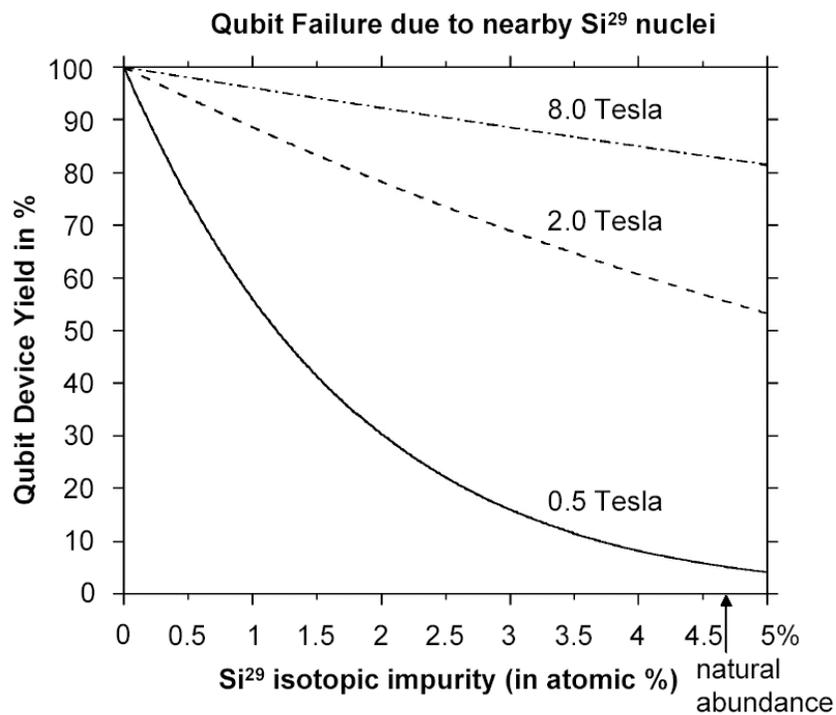
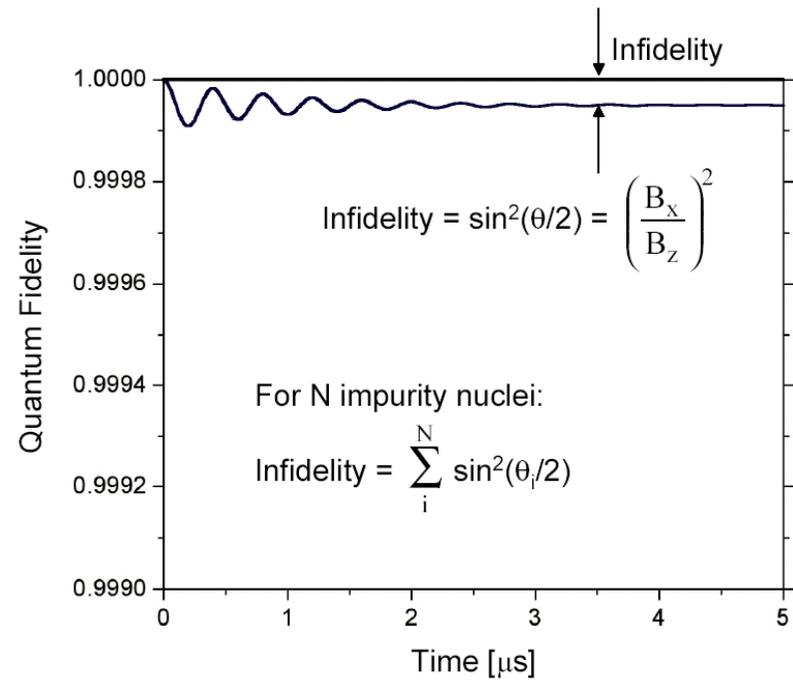
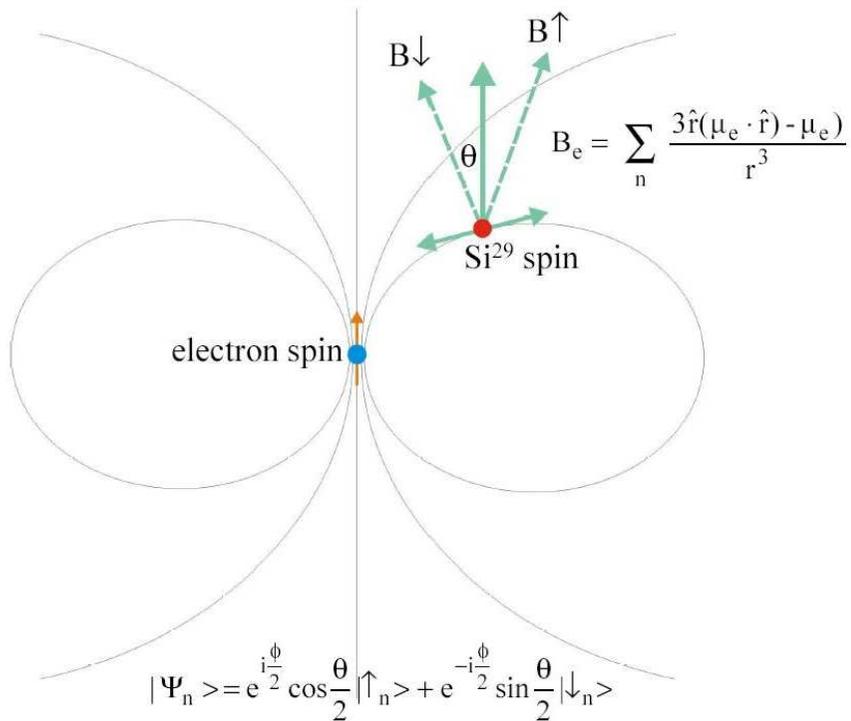


$$\Delta\omega = \sum_n \frac{1}{\hbar r^3} [3(\mu_n \cdot \hat{r})(\mu_e \cdot \hat{r}) - (\mu_n \cdot \mu_e)]$$

$$e^{i \int_0^t \Delta\omega(t') dt'} = e^{i \int_0^t A \sqrt{r'} dt'} = e^{i \frac{2}{3} A t^{3/2}} \Rightarrow \exp\{- (t/t_0)^3\}$$

Spin Dephasing Lifetime $t_0 \sim \left[\frac{1}{N_{Si^{29}}} \right]^\alpha$ model dependent exponent

Isotopic impurity density \uparrow



Nuclear Spins can also produce an initial step Projection Error in many types of macro-scopic qubits!

A macroscopic qubit can cause un-intentional off-resonant NMR, whenever electric currents exist.

For example, a flux quantum is associated with $B \sim 1$ milli-gauss

This can produce an NMR frequency of $f \sim 6\text{Hz}$

At a clock period of 100nsec,

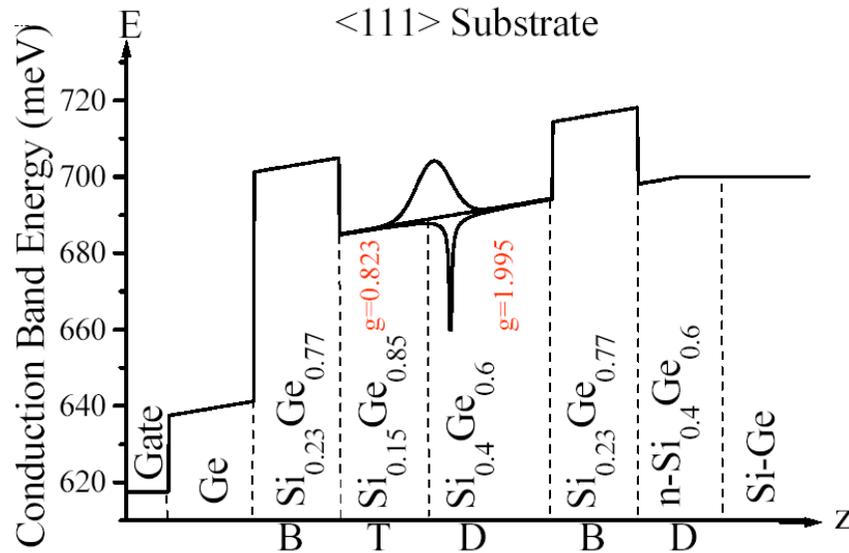
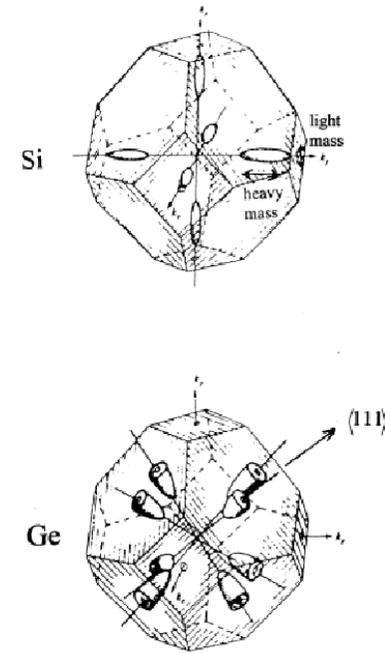
this can disturb the nuclear spins by $\sim 10^{-7}$ radians

This would produce an infidelity of $\sim 10^{-14}$

But if there are $\sim 10^{10}$ nuclear spins in the sample,

The total infidelity would be $\sim 10^{-4}$, enough to be concerned.

Conduction bands in Si & Ge



Band Structure Engineering is famous for controlling effective mass.

Thus it controls g-factor also:

$$g-2 \sim -\frac{\text{spin-orbit coupling}}{E_g} \left(\frac{m}{m_{\text{eff}}} - 1 \right)$$

“g-factor engineering”

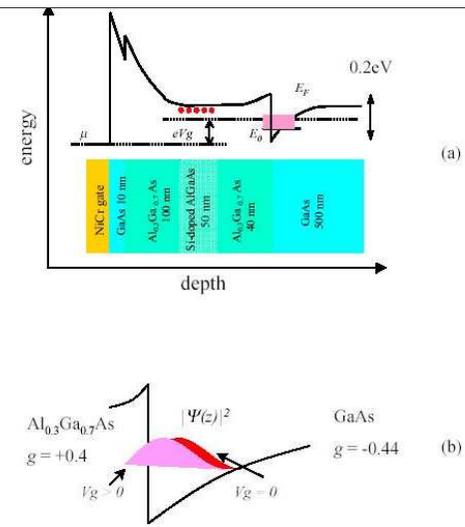
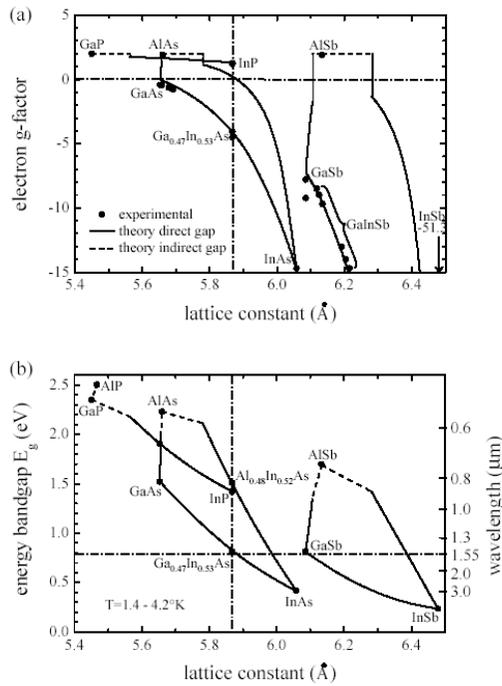


Figure 1. (a) Schematic diagram of the conduction band of the modulation doped GaAs/Al_{0.3}Ga_{0.7}As heterostructure. Fermi level, E_F , the chemical potential, μ , and the lowest energy level of the quantum well, E_0 , are indicated. (b) The two-dimensional electrons are trapped in the "triangle" shaped quantum well near the interface of the GaAs and Al_{0.3}Ga_{0.7}As materials. The electron wavefunction shifts back and forth for bias voltage $V_g = 0$ and a positive bias voltage.

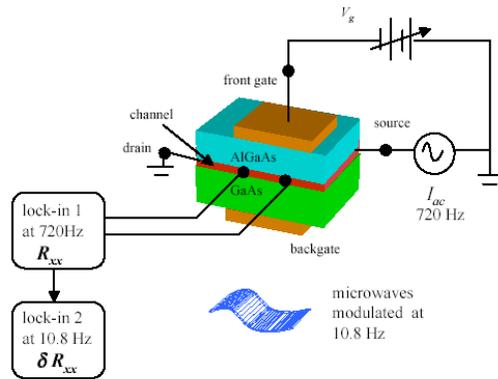
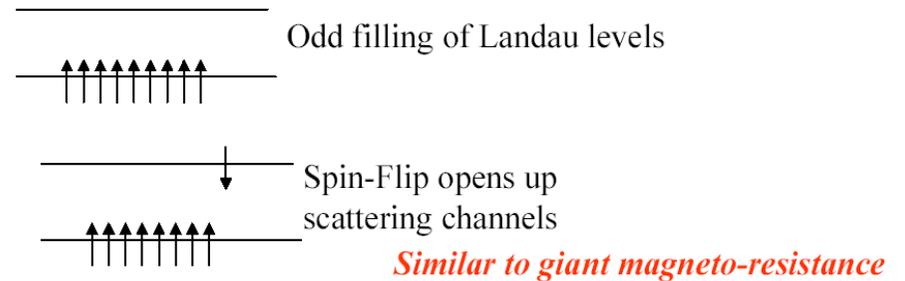


Figure 2. Diagram of the experimental setup for monitoring electron spin resonance and for controlling the spin orientation.

Spin Detection Method:



Single Spin Detection Methods:

1. Spin Dependent tunneling from a quantum well (change charge state from $1 \Rightarrow 0$)
2. Form singlet bound state versus unbound triplet (change charge state from $1 \Rightarrow 2$)

Then detect single charge in an FET or Quantum Point Contact transistor

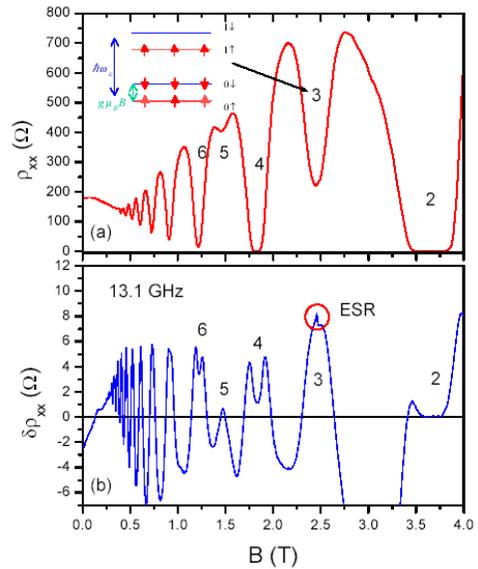


Figure 3. (a) Typical trace of the resistivity ρ_{xx} as a function of the magnetic field. Landau level filling factors ν are indicated. Inset: energy diagram for the case of $\nu = 3$. (b) The microwave radiation induced resistivity change $\delta\rho_{xx}$. Note the ESR feature around $\nu = 3$.

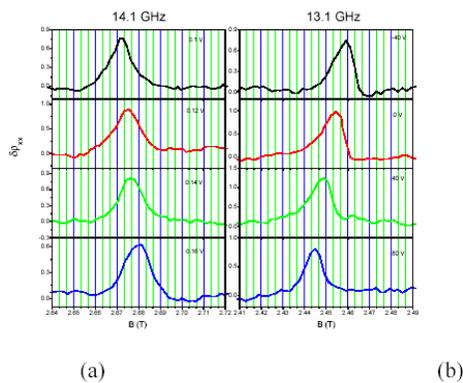
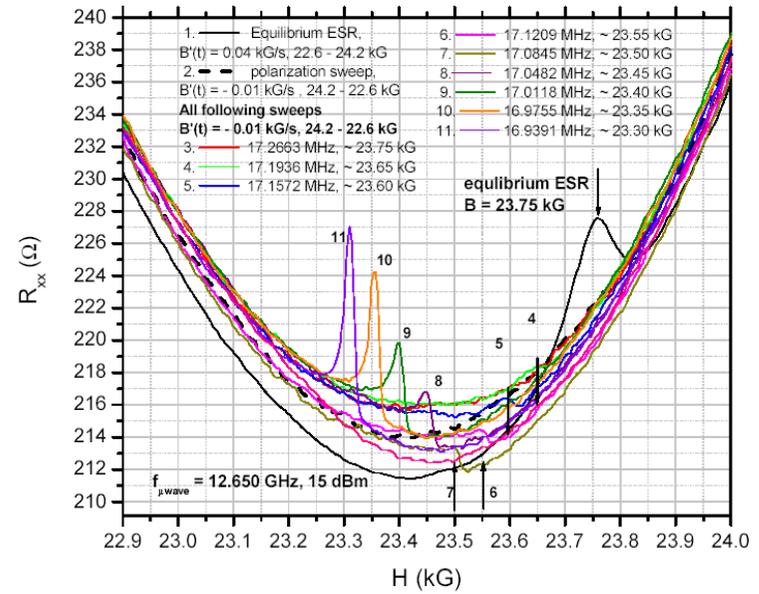


Figure 4. Electrically detected electron spin resonance spectra at a sequence of gate voltages for (a), a front gate and, (b), a back gate. For the front gate case, the resonant peak moves progressively to higher magnetic field when the amplitude of the bias gate voltage is increased. In contrast, the peak shifts towards lower field for the back gate case. (The non-resonant background of the signal was subtracted for clarity. This task can be easily done since the background varies smoothly over a much broader range of magnetic field.)

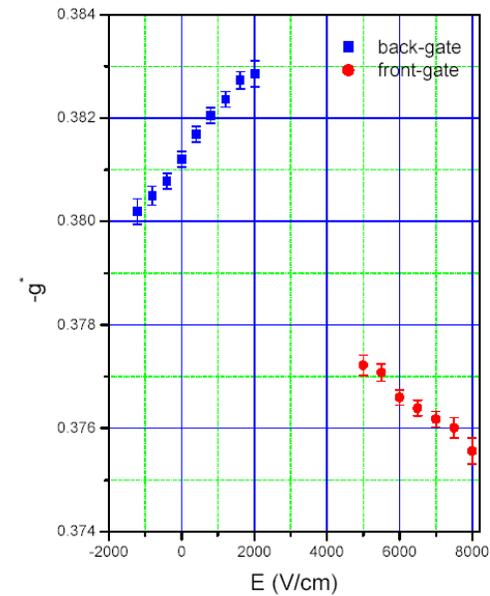
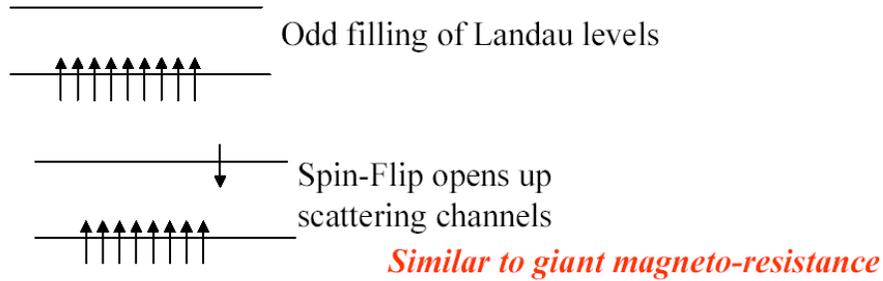


Figure 5. Experimentally determined electronic g-factor as a function of the applied electric field for both the front and back gate. The electric field is simply the applied voltage divided by insulator thickness. In this plot, no attempt was made to include space charge self-consistently.

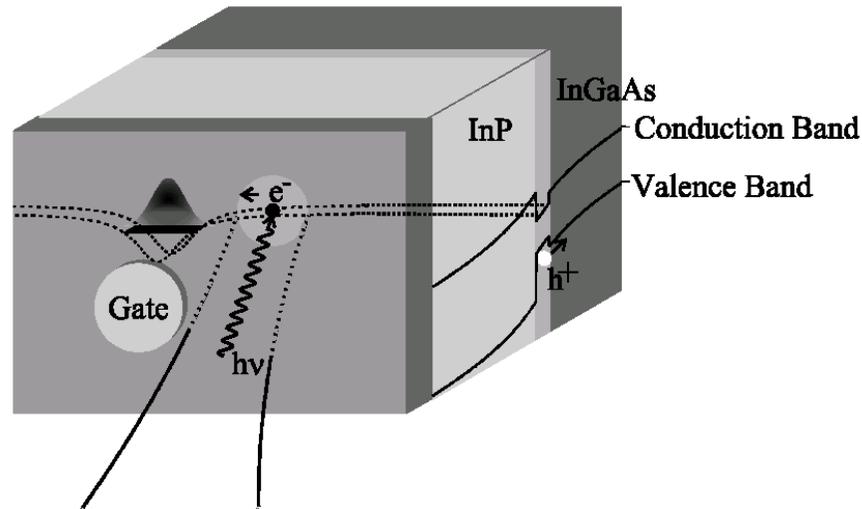
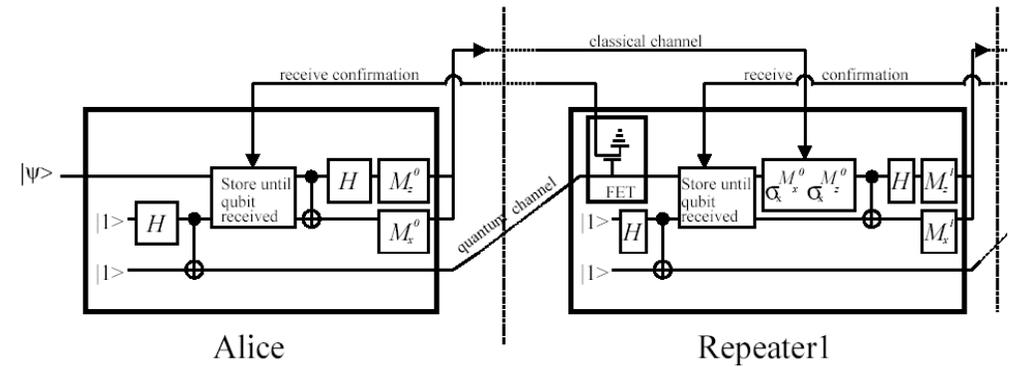
Spin Detection Method:



Single Spin Detection Methods:

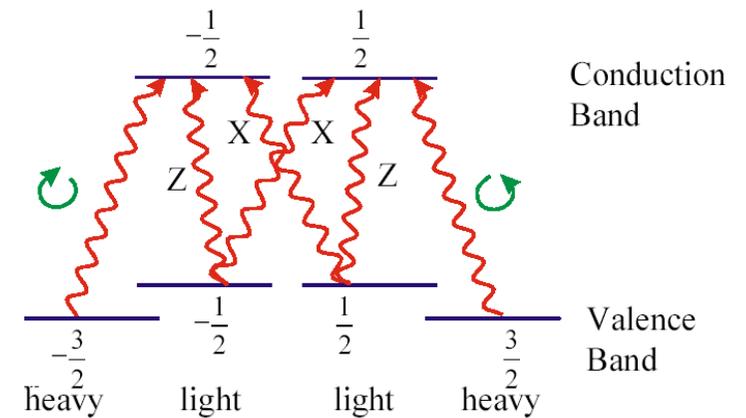
1. Spin Dependent tunneling from a quantum well
(change charge state from $1 \Rightarrow 0$)
2. Form singlet bound state versus unbound triplet
(change charge state from $1 \Rightarrow 2$)

Then detect single charge in an FET or Quantum Point Contact transistor



Entanglement Preserving Spin-Coherent Semiconductor Photodetector

Entanglement Preserving Photo-Detector



The hole wave functions in Semiconductors:

$$|J=3/2, L=1, S=1/2\rangle$$

heavy $|m_j = 3/2\rangle = |m_l = 1, m_s = 1/2\rangle$

light $|m_j = 1/2\rangle = |m_l = 1, m_s = -1/2\rangle + |m_l = 0, m_s = 1/2\rangle$

light $|m_j = -1/2\rangle = |m_l = -1, m_s = 1/2\rangle - |m_l = 0, m_s = -1/2\rangle$

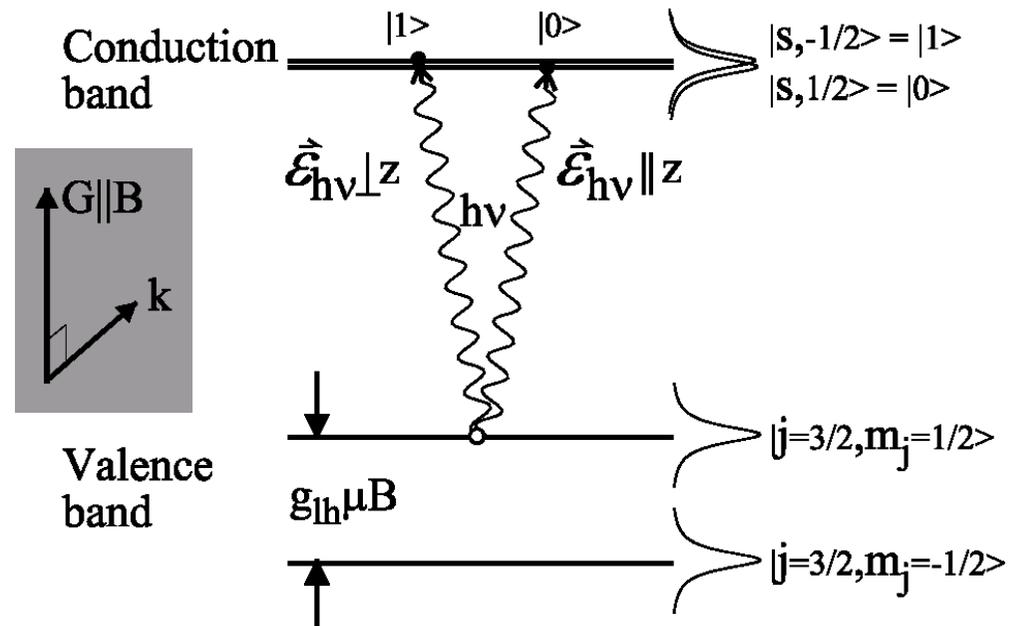
heavy $|m_j = -3/2\rangle = |m_l = -1, m_s = -1/2\rangle$

The electron wave function in III-V semiconductors:

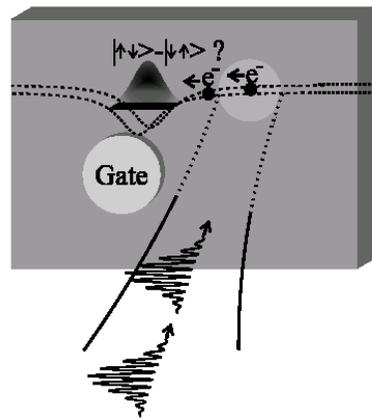
$$|L=0, S=1/2\rangle$$

electron $|m_l = 0, m_s = 1/2\rangle$

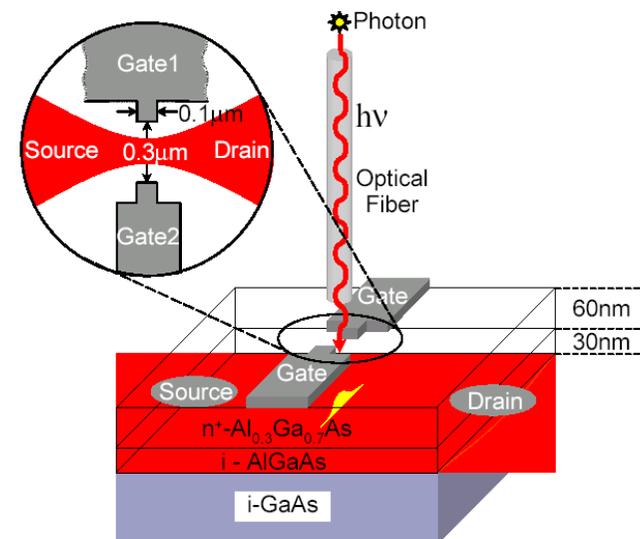
electron $|m_l = 0, m_s = -1/2\rangle$



Detecting Entanglement:



Parametric Decay Photon Pair Generator



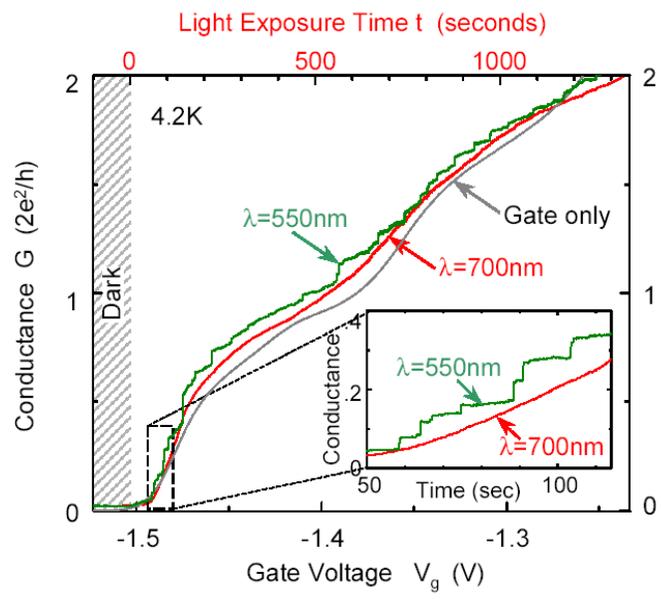
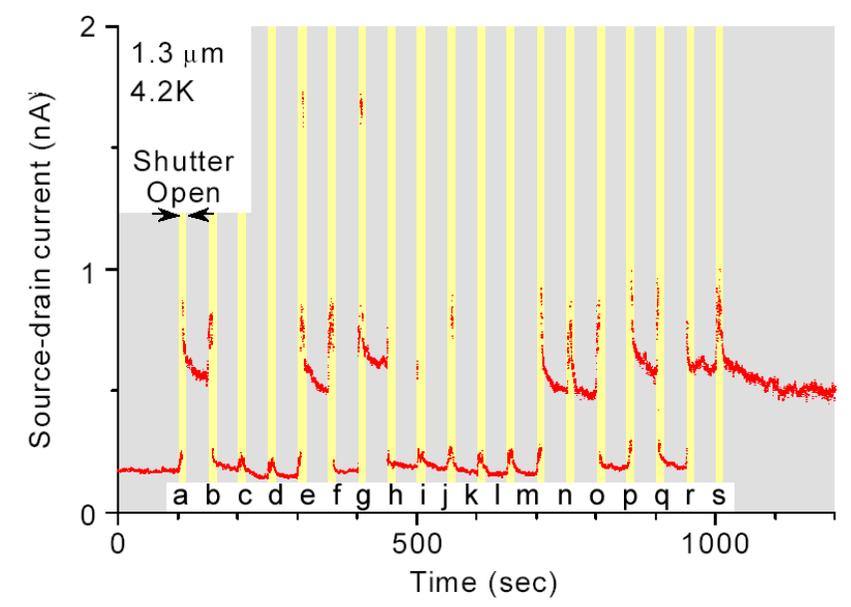
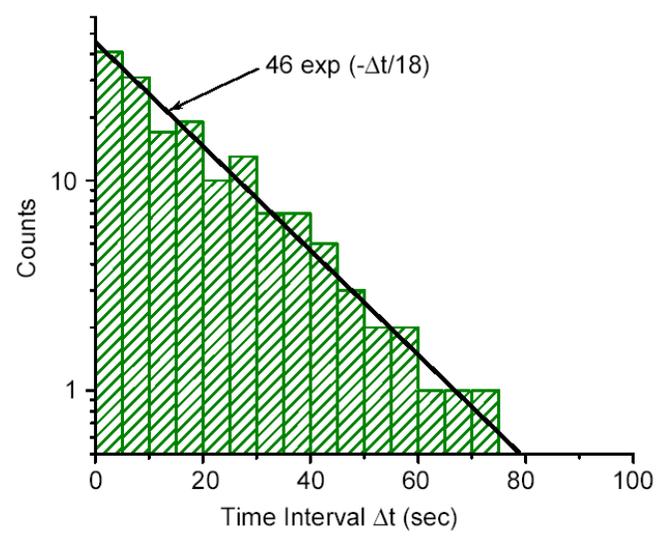
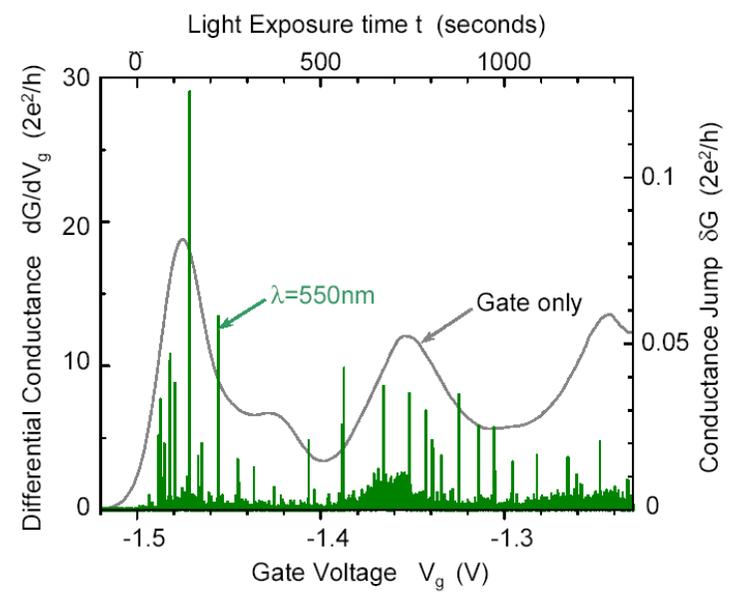


Fig. 1(a) Hideo Kosaka et al.



Hideo Kosaka Fig. 5

Direct Observation of the Precession of Individual Paramagnetic Spins on Oxidized Silicon Surfaces

Y. Manassen, R. J. Hamers, J. E. Demuth, and A. J. Castellano, Jr.
 IBM Research Division, T. J. Watson Research Center, Yorktown Heights, New York 10598
 (Received 12 December 1988)

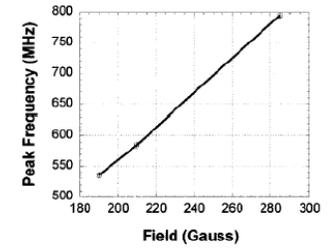
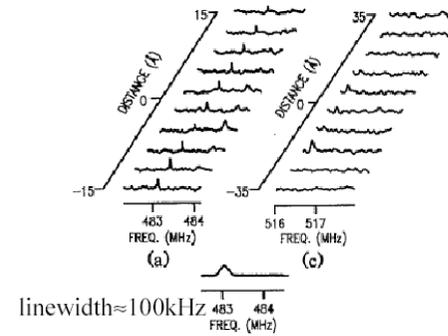


FIG. 5. Plot of the center frequency of STM-ESR peaks on clusters as a function of the applied magnetic field. From this, we obtain a value of $g=2\pm 0.1$.

linewidth ≈ 100 kHz

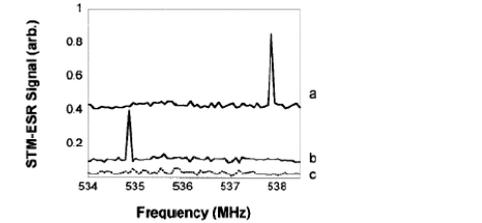


FIG. 3. STM-ESR spectra of (a), (b) two different areas (a few nm apart) of the molecule-covered sample and (c) bare HOPG. The graphs are shifted vertically for clarity.

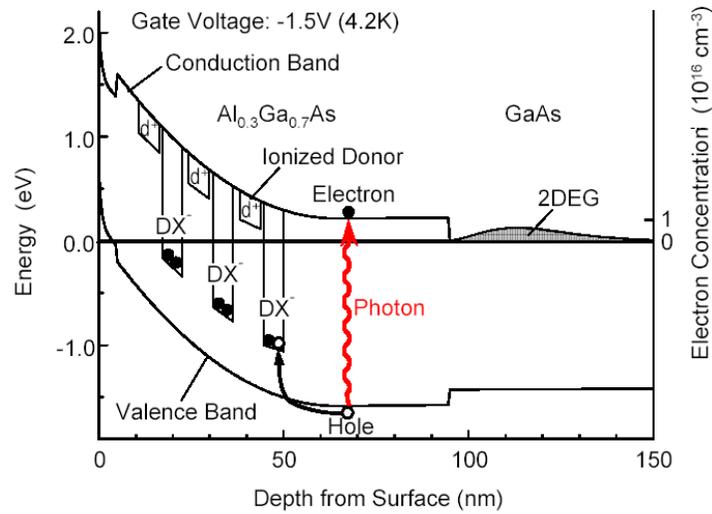
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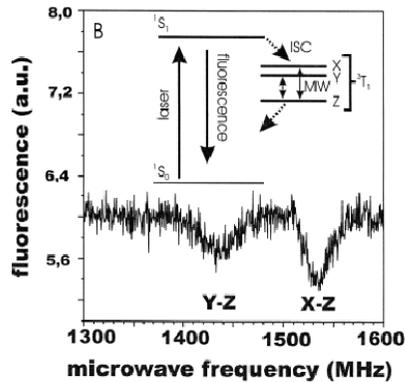
21 JANUARY 2002

Electronic spin detection in molecules using scanning-tunneling-microscopy-assisted electron-spin resonance

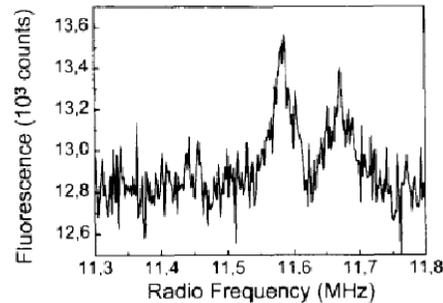
C. Durkan¹⁾ and M. E. Welland
 Nanoscale Science Laboratory, Department of Engineering, University of Cambridge, Trumpington Street, Cambridge CB2 3PE, United Kingdom



Pentacene in *p*-terphenyl



ESR

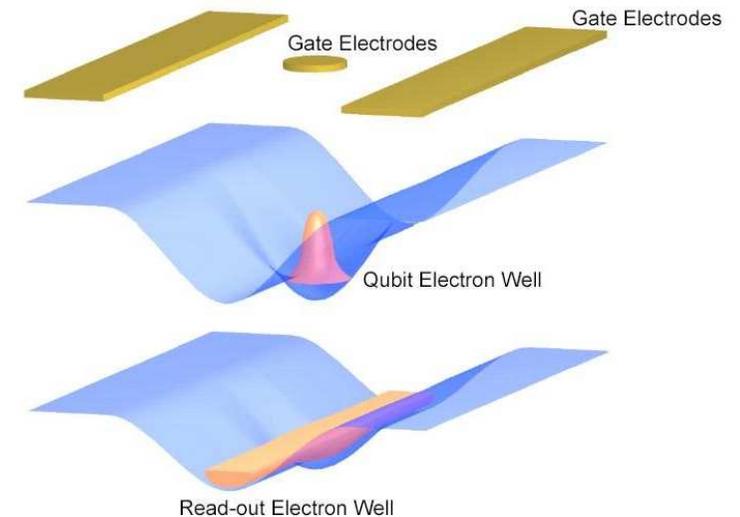


NMR

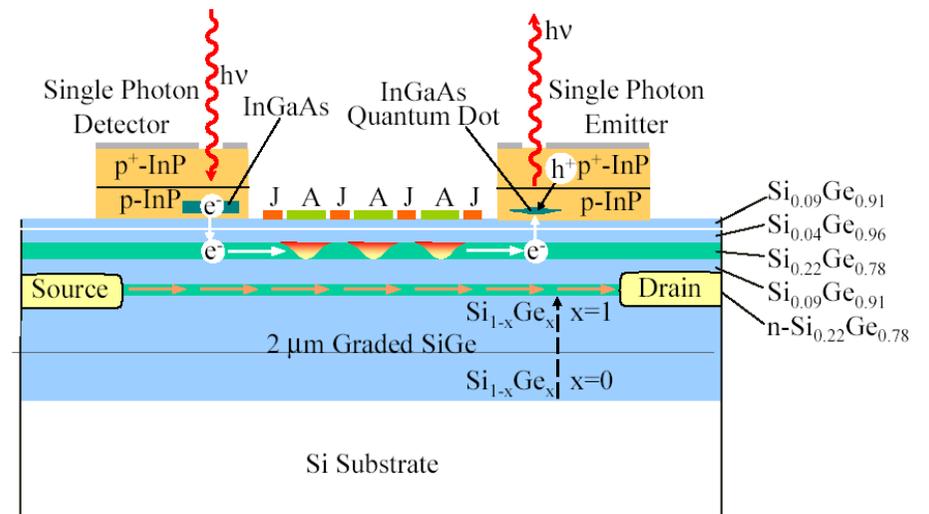
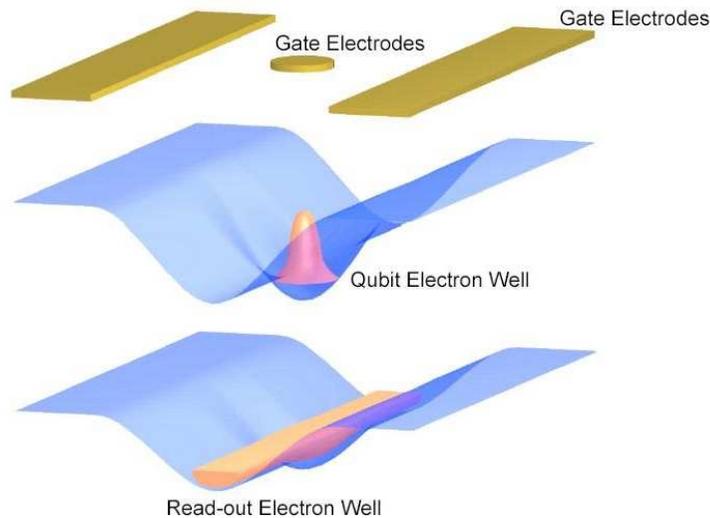
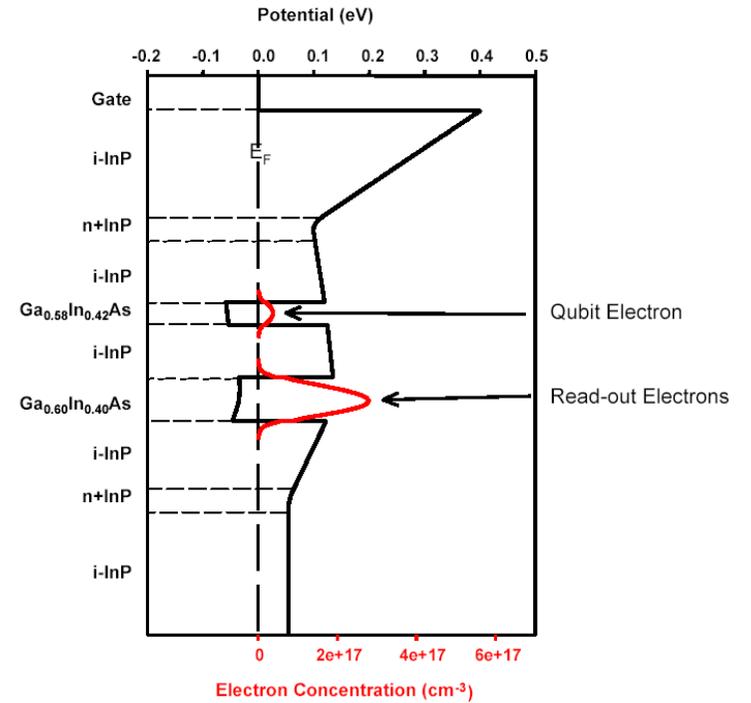
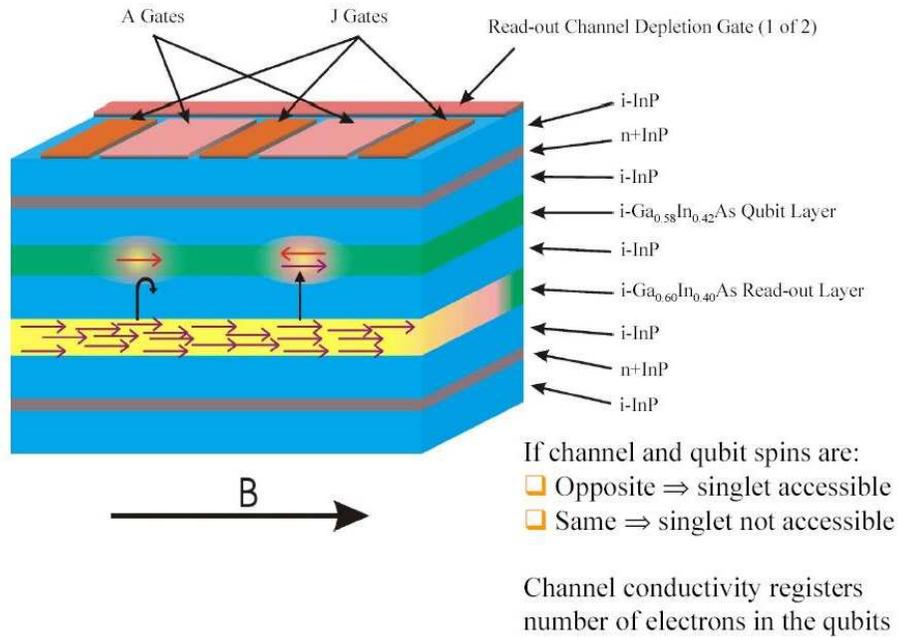
Fig. 3. ENDOR spectrum of a single molecule recorded with a microwave excitation at 1512 MHz, corresponding to the central peak in the ESR spectrum in Fig. 2. 3500 scans with a duration of 0.8 s have been accumulated.

J. Wrachtrup, A. Gruber, *TU Chemnitz, Inst. of Physics, 09107 Chemnitz, Germany*
 SSNMR vol.11, 59-64 (1998)

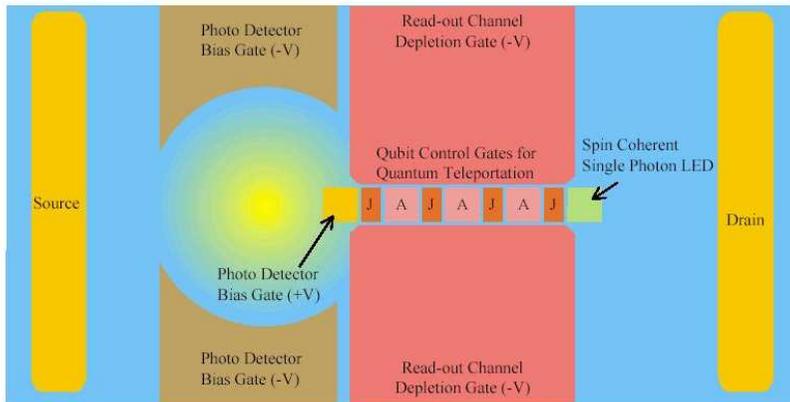
J. Wrachtrup, A. Gruber, *TU Chemnitz, Inst. of Physics, 09107 Chemnitz, Germany*
 Chem. Phys. Lett. 267, 179-185 (1997)



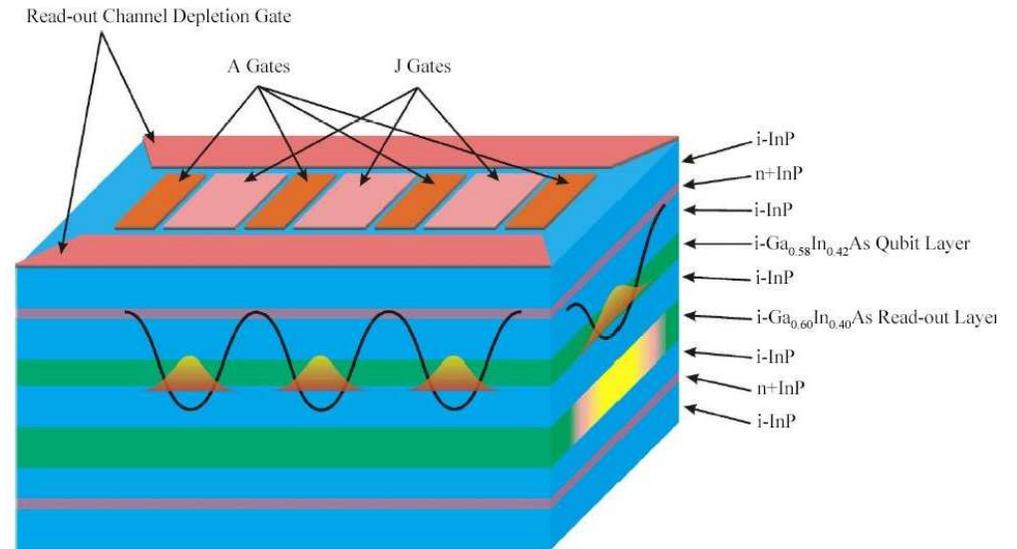
Qubit Read-out Scheme



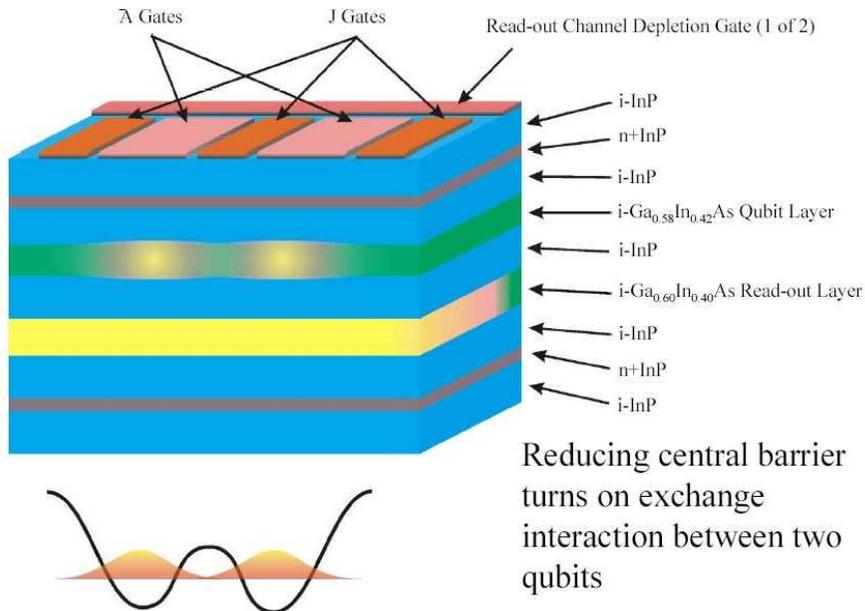
Teleportation Repeater Top-view Schematic



Gates for Quantum Repeater



Two qubit interaction



Conclusions:

1. There is *little or no penalty* in using electron spins in Silicon versus nuclear spins.
2. Electron spin Dephasing is already slow enough to easily permit error correction.
3. Single photon and single photo-electric charge detection is readily performed in transistor-like structures.
4. The next big step prior to an electron spin CNOT gate is the readout of a single electron spin orientation.
5. In the long run;
Carbon-based electron-spin hosts may allow room temperature operation of a quantum information processor.

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Single photon detection by Photoconductive Gain experiment
- Dr. Hans Robinson, post-doc, UCLA EE Dept.
Spin readout device design
- Dr. Rutger Vrijen, former post-doc, now at SUN Microsystems Lab
Entanglement preserving photodetector

QUANTUM INFORMATION SPINTRONICS

Electron Spins logic & storage
(in Si)

Nuclear Spins storage
(in Si)

Electron Spins tele-communication
(in III-V)

...requires confinement energy $> kT$,
thus small quantum dots are not needed;
e-beam lithography is adequate.

Technology Generation	Technology	Gate Count
1	Quantum Key Distribution (fairly Robust) - Heralded –Single-Photon transmitters (Yamamoto et al)	0
2	Quantum Repeater – long distance EPR pair distribution detect photo electron (don't perturb spin)	3
3	Quantum memory – higher secure throughput	≥ 10
4	Quantum Error Correction	≥ 40
5	Quantum Factorization Engine	≥ 1000