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Cosmology and Structure Formation

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Probing cosmic acceleration
with galaxy clusters



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Outline

- **Motivation**
- **Methodology**
 - Known Problems (?)
 - Possible Solutions (?)
- **Applications**
 - Dark Energy Effects
 - Model Comparison
- **Conclusions**

Motivation



Why Cluster ?

: The largest virialized objects in the Univ

Formation solely depends on gravity irrelevant to gas dynamics, SF, feedback

Abundance and evolution $\longleftrightarrow P_{\text{lin}}(k, a)$

Thus, abundance and distribution are determined by geometry of Univ and power spectrum

$$dn(M, z) = \frac{\rho_m^0}{M} \frac{d \ln \sigma(M, z)}{dM} \boxed{f(M)} dM$$

→ cosmology
→ simulation

$$\sigma_R^2(a) \equiv \left\langle \left| \frac{\delta M}{M(R, a)} \right|^2 \right\rangle = \frac{1}{2\pi^2} \int_0^\infty k^2 \boxed{P(k, a)} |W(kR)|^2 dk, \quad f(M) : \text{mass function.}$$

Methodology - I

- Problems (?)
 - Press-Schechter(PS) predicts too few high mass clusters and too many low mass ones
 - Too simple : not realistic ?
- Solutions (?) : traditional way
 - Find proper $f(M)$ based on Simulation
- Alternative
 - Check fundamental quantities
 σ, δ_c

- Evolutions outside and inside the overdensity region (R)

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}(\rho_m + \rho_{de}) = \frac{8\pi G}{3}\rho_{cr}$$

$$\frac{\ddot{R}}{R} = -\frac{4\pi G}{3}[\rho_{cluster} + (1 + 3\omega_{de})\rho_{dec}]$$

$$\dot{\rho}_{dec} + 3(1 + \omega_{de})\left(\frac{\dot{R}}{R}\right)\rho_{dec} = \alpha\Gamma,$$

$$\text{where } \Gamma = 3(1 + \omega_{de})\left(\frac{\dot{R}}{R} - \frac{\dot{a}}{a}\right)\rho_{dec} \quad \text{with } 0 \leq \alpha \leq 1$$

- Find (semi)-analytic solutions for a R of spherical collapse model for general DE : compared to EdS model solution (cycloid) of Peebles

Methodology - II

- EdS cycloid nice but without DE solution

$$\frac{d^2 R}{dt^2} = -\frac{GM}{R^2} = -\frac{4\pi G}{3} \langle \rho \rangle (1 + \langle \delta \rangle) R$$

$$\frac{R}{R_m} = \frac{1}{2}(1 - \cos \eta) \quad \frac{t}{t_m} = \frac{1}{\pi}(\eta - \sin \eta)$$

- General solutions with DE (SL 10, SL & Ng 10)

$$x = \frac{a}{a_{ta}}, \quad y = \frac{R}{R_{ta}}$$

$$d\tau \equiv H_{ta} \sqrt{\Omega_m(x_{ta})} dt, \quad \zeta \equiv \frac{\rho_{cluster}}{\rho_m}|_{z_{ta}}, \quad Q_{ta} \equiv \frac{\rho_m}{\rho_{de}}|_{z_{ta}} = \frac{\Omega_m^0}{\Omega_{de}^0} (1 + z_{ta})^{-3\omega_{de}}$$

$$\frac{2}{3} x^{\frac{3}{2}} F\left[\frac{1}{2}, -\frac{1}{2\omega_{de}}, 1 - \frac{1}{2\omega_{de}}, -\frac{x^{-3\omega_{de}}}{Q_{ta}}\right] = \tau$$

$$\zeta_{13} = \left(\frac{\pi/2}{\frac{2}{3} F\left[\frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \frac{-1}{Q_{ta13}}\right]} \right)^2 = \left(\frac{3\pi}{4} \right)^2 \left(F\left[\frac{1}{2}, \frac{3}{2}, \frac{5}{2}, -\frac{(1 - \Omega_m^0)}{\Omega_m^0} (1 + z_{ta})^{-1}\right] \right)^{-2}$$

- General solutions (continued)

$$\zeta_{sk} = \left(\frac{3\pi}{4} \right)^2 \Omega_{mta}^{-0.724 + 0.157 \Omega_{mta} + \alpha(1 + \omega_{de})(1 + 3\omega_{de})(0.064 - 0.368 \Omega_{mta})}$$

$$\begin{aligned} & \sqrt{y(1-y)} - \text{ArcSin}[\sqrt{y}] + \frac{\pi}{2} + \left(-\frac{(1+3\omega_{de})}{3} \Omega_{mta} \right)^{2.7+0.1\Omega_{mta}-3.4(1+\omega_{de})-0.02\zeta_{sk}+3(z_{ta}-0.6)} \\ &= \sqrt{\zeta_{sk}}(\tau - \tau_{ta}) \\ &= \frac{2}{3} \sqrt{\zeta_{sk}} \left(x^{\frac{3}{2}} F\left[-\frac{1}{2\omega_{de}}, \frac{1}{2}, 1 - \frac{1}{2\omega_{de}}, -\frac{x^{-3\omega_{de}}}{Q_{ta}}\right] - F\left[-\frac{1}{2\omega_{de}}, \frac{1}{2}, 1 - \frac{1}{2\omega_{de}}, -\frac{1}{Q_{ta}}\right] \right). \end{aligned}$$

- EdS virial solution instead of collapsed one by using $z_{vir} = 1/2 z_{ta}$

$$\frac{2}{3} x^{\frac{3}{2}} = \tau \quad \zeta = \left(\frac{3\pi}{4} \right)^2 \quad \frac{3}{20} (12\pi)^{\frac{2}{3}} \simeq \boxed{1.69}$$

$$\sqrt{y(1-y)} - \text{ArcSin}[\sqrt{y}] + \frac{\pi}{2} = \sqrt{\zeta}(\tau - \frac{2}{3}) = \frac{\pi}{2} (x^{\frac{3}{2}} - 1)$$

$$\Delta_c^{\text{EdS}} = \zeta \left(\frac{x_c}{y_{vir}} \right)^3 = \left(\frac{3\pi}{4} \right)^2 \left(\frac{2^{\frac{2}{3}}}{2^{-1}} \right)^3 = 18\pi^2 \simeq \boxed{178}$$

$$\Delta_{vir}^{\text{EdS}} = \zeta \left(\frac{x_{vir}}{y_{vir}} \right)^3 = \left(\frac{3\pi}{4} \right)^2 \left(\frac{(\frac{1}{\pi} + \frac{3}{2})^{\frac{2}{3}}}{2^{-1}} \right)^3 = 18\pi^2 \left(\frac{1}{2\pi} + \frac{3}{4} \right)^2 \simeq \boxed{147}$$

$$\delta_{lin}(z_{vir}) = \frac{3}{5} (\sqrt{\zeta})^{\frac{2}{3}} \left(\left(\frac{3}{4} + \frac{9\pi}{8} \right) \frac{1}{\sqrt{\zeta}} \right)^{\frac{2}{3}} = \frac{3}{20} (6 + 9\pi)^{\frac{2}{3}} \simeq \boxed{1.58}$$

Application - I

- Do we need to care about small difference between 1.69 and 1.58 ?

- PS mass function includes δ_c

- PS used in $dn(M,z)$ $f_{PS}(M,z) = \sqrt{\frac{2}{\pi}} \frac{\delta_c}{\sigma} \exp\left[-\frac{\delta_c^2}{2\sigma^2}\right]$

- Nonlinear overdensity in simulation use 200 in any case instead of 178, thus, is 147 important ?

- However, DE also alters P_{lin} and thus σ

- DE is usually not clustered at small scale due to relativistic dispersion of its fluctuation and **usually** DE effects on σ_8 is ignored

- DE effects on P_{lin} , σ (Ma et al 99, SL & Ng 10))

$$P(k, a) = A_Q k^{n_s} T_Q^2(k) \left(\frac{D(a)}{D(a_0)} \right)^2$$

$$A_Q = 2\pi^2 \delta_H^2 (c/H_0)^{n_s+3}$$

$$\begin{aligned} \delta_H &= 2.05 \times 10^{-5} \alpha_Q^{-1} (\Omega_m)^{c_1+c_2 \ln \Omega_m} \exp[c_3(n_s-1) + c_4(n_s-1)^2], \text{ with} \\ c_1 &= -0.789|\omega_Q|^{0.0754-0.211 \ln |\omega_Q|}, c_2 = -0.118 - 0.0727\omega_Q, c_3 = -1.037, c_4 = -0.138 \\ \alpha &= (-\omega_Q)^s \text{ with } s = (0.012 - 0.036\omega_Q - 0.017\omega_Q^{-1})(1 - \Omega_m(a)) \\ &\quad + (0.098 + 0.029\omega_Q - 0.085\omega_Q^{-1}) \ln \Omega_m(a). \end{aligned}$$

$$\sigma_R^2(a) \equiv \left\langle \left| \frac{\delta M}{M(R, a)} \right|^2 \right\rangle = \frac{1}{2\pi^2} \int_0^\infty k^2 P(k, a) |W(kR)|^2 dk.$$

$$\sigma_8(\omega_Q) = (-\omega_Q)^{0.72+0.36\omega_Q} \sigma_8(\omega_Q = -1)$$

$$\sigma(M, z) \simeq (-\omega_Q)^{0.72+0.36\omega_Q} \left(3.90 - 0.215 \log \left[\frac{M}{h^{-1} M_\odot} \right] \right) \left(\frac{D_g(z)}{D_g(z_0)} \right)$$

$$R_8 = 8 h^{-1} \text{Mpc as } M_8 = 5.95 \times 10^{14} \Omega_m^0 h^{-1} M_\odot$$

Application - II

- DE effects on Matter power spectrum and σ_8

More matter in the past for the smaller value of ω to give larger value of δ at present

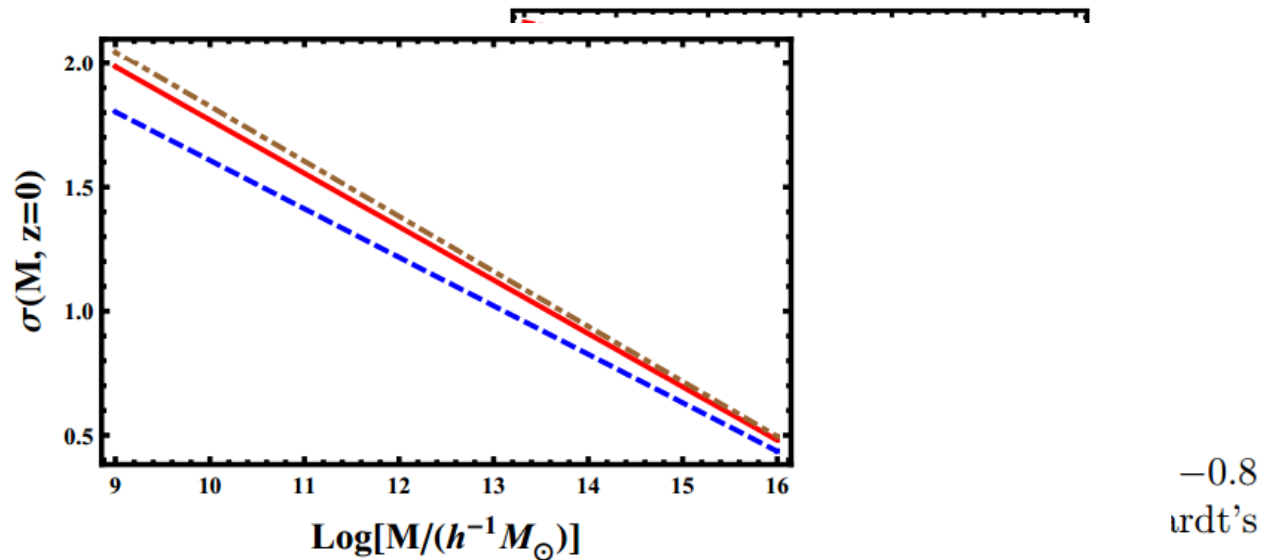


Figure 3: $\sigma(M, z = 0)$ for different values of $\omega_Q = -1.1, -1.0$, and -0.8 (from top to bottom)

- DE effects on $\sigma(M, z)$

Application - III

- DE effects on Growth factor and Volume (SL & Ng 10)

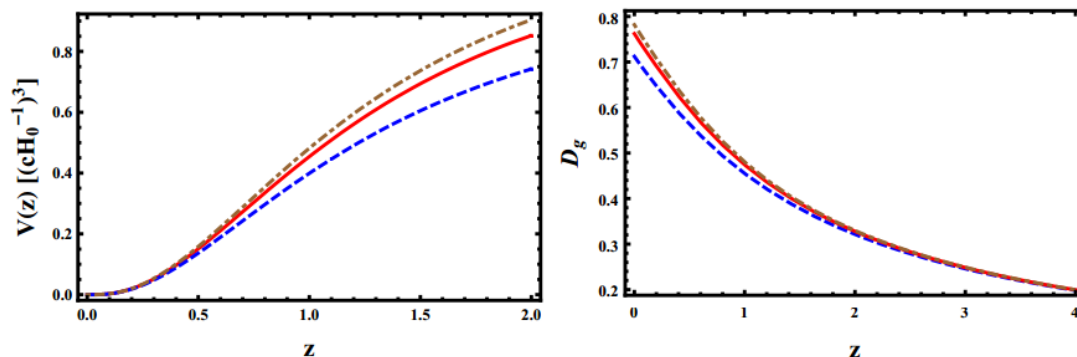


Figure 4: a) $V(z)$ for different values of $\omega_Q = -1.1, -1.0$, and -0.8 (from top to bottom). b) $D_g(z)$ for the same values of ω_Q as in the left panel.

- comoving number density of virialized objects

$$dn(M, z) = \sqrt{\frac{2}{\pi}} \frac{\rho_m^0}{M^2} \left| \frac{d \ln \sigma}{d \ln M} \right| \frac{\delta_c}{\sigma} \exp \left[-\frac{\delta_c^2}{2\sigma^2} \right] dM$$

$$f_{\text{ST}}(\sigma) = A \sqrt{\frac{2b}{\pi}} \exp \left[-\frac{b\delta_c^2}{2\sigma^2} \right] \left[1 + \left(\frac{\sigma^2}{b\delta_c^2} \right)^p \right] \frac{\delta_c}{\sigma}$$

$$f_{\text{mod}}(\sigma, z) = \tilde{A} \sqrt{\frac{2}{\pi}} \exp \left[-\frac{\tilde{b}\delta_c^2}{2\sigma^2} \right] \left[1 + \left(\frac{\sigma^2}{\tilde{b}\delta_c^2} \right)^{\tilde{p}} \right] \left(\frac{\delta_c \sqrt{\tilde{b}}}{\sigma} \right)^{\tilde{q}}$$

Application - IV

- Results (data from Calberg et al 96)

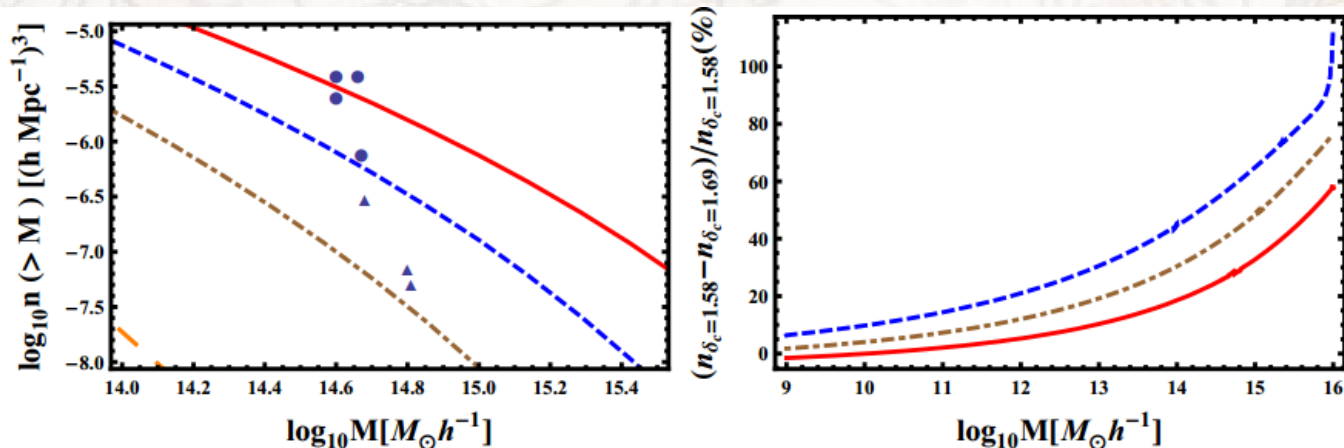


Figure 5: a) The comoving number density of clusters n of mass greater than M for different values of $z = 0, 0.5, 1.0$, and 2.0 (from top to bottom) when $\omega_Q = -1$ and $\delta_c = 1.58$. The circular ($z \simeq 0$) and triangular ($0.18 \leq z \leq 0.85$) dots represent the data from Ref. [33]. b) Errors of n when we use the correct threshold density contrast $\delta_c = 1.58$ instead of 1.69 for different values of $z = 0, 0.5$, and 1.0 (from bottom to top).

Conclusions

- We found the correct values of δ_c and Δ_{lin}
- Dark energy is known to be clustered only at large scale. However, it still affects on not only V and D but also P_{lin} and thus σ_8
- We apply correct values of δ_c and effects of DE on σ_8 to solve shortage problem of high mass clusters with PS formalism
- Low mass is still problem with PS. Not only gravity?